

Investigating moiré interlayer excitons under the influence of atomic reconstructions

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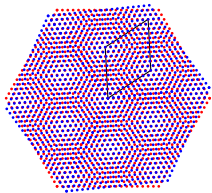
University of Bremen

Institute for Theoretical Physics

18.03.2024

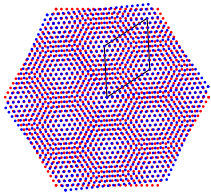


Motivation

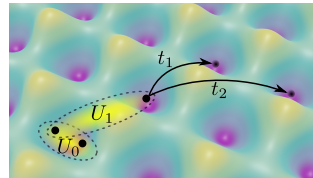


Twisted TMDs

Motivation

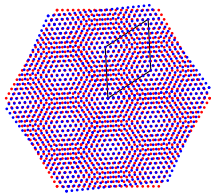


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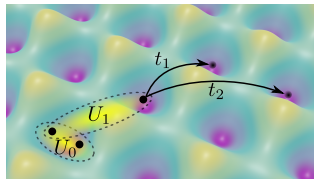


Götting *et al.*, Phys. Rev. B
105, 165419 (2022)

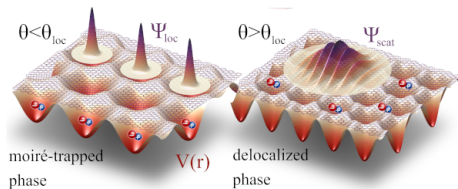
Motivation



Twisted TMDs



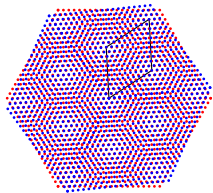
Götting *et al.*, Phys. Rev. B
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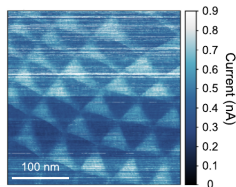
Quantum Phases

Brem *et al.* Nano Lett., 20, 12, 8534-8540
(2020)

Motivation



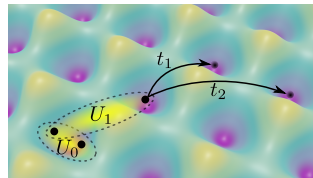
Twisted TMDs



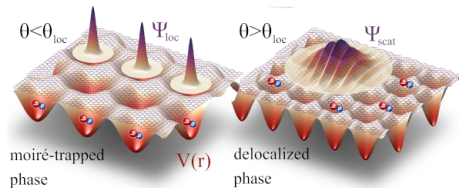
Atomic
Reconstructions

Rosenberger *et al.*
ACS Nano, **14**, 4,
4550-4558 (2020)

Quantum Phases



Götting *et al.*, Phys. Rev. B
105, 165419 (2022)

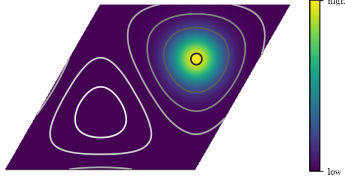


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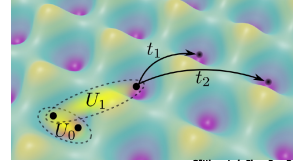
Outline



Interlayer
Excitons

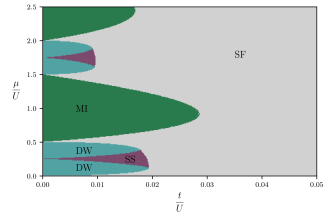


Mapping onto a
Bose-Hubbard Model

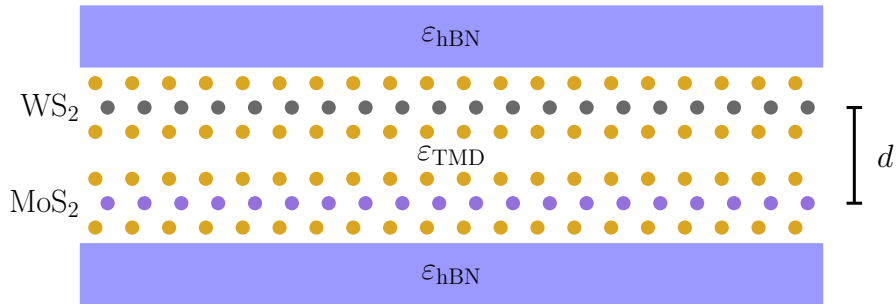


Götting et al., Phys. Rev. B
105, 165419 (2022)

Influence of Atomic
Reconstructions on the
Quantum Phases of IXs

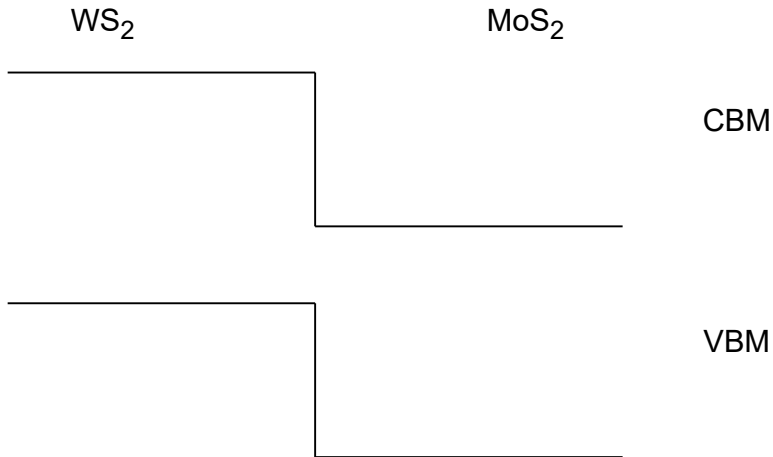


Stacking two TMDs

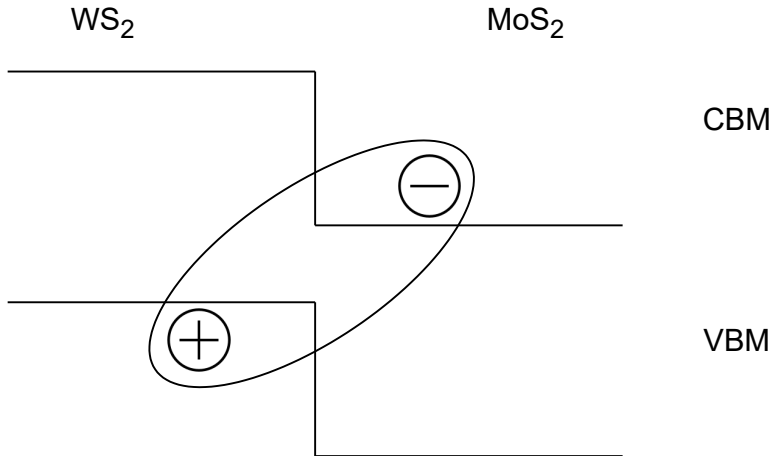


Stacking structure: R_X^X ,
X: chalcogen atom

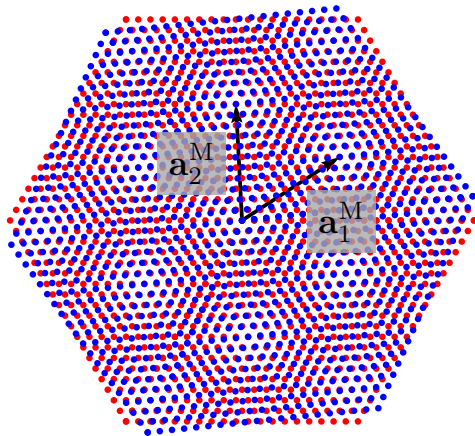
Type-II Band Alignment

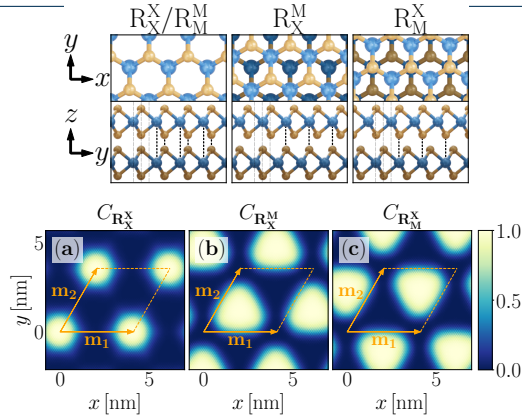


Exciton Formation



Introducing Twist Angle

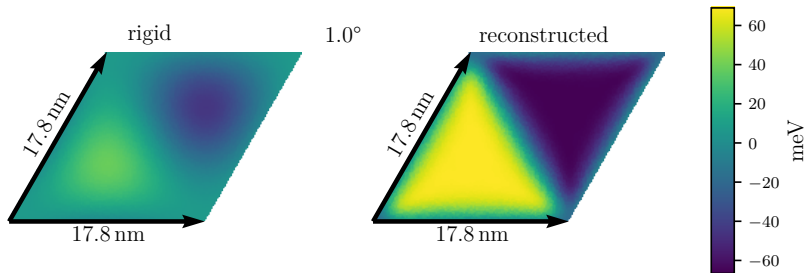




AG Bester, Universität Hamburg

- band gap as effective potential for excitons:

$$V^M(\mathbf{r}) = \sum_{\mathbf{G}} V_{\mathbf{G}} e^{i\mathbf{G}\mathbf{r}}$$

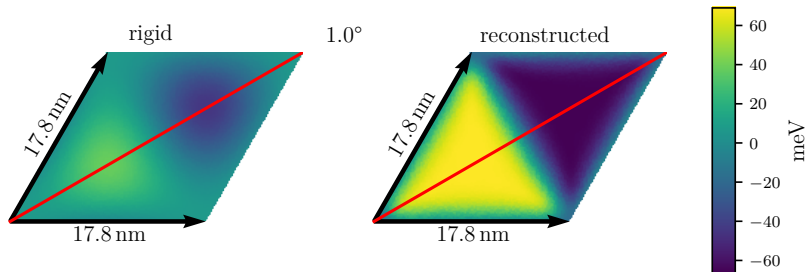


Wu *et al.*, Phys. Rev. Lett. 118, 1474019 (2017)

Nielsen *et al.*, Phys. Rev. B 108, 045402 (2023)

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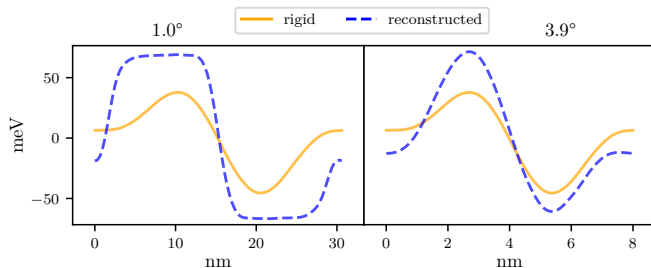


Wu *et al.*, Phys. Rev. Lett. 118, 1474019 (2017)

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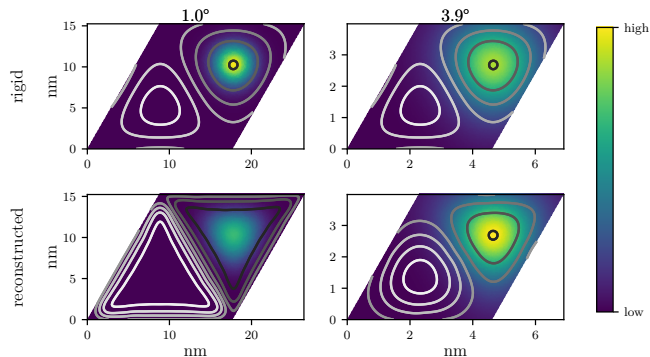
IX Hamiltonian

$$H = -\frac{\hbar^2}{2M}\Delta + V^{\text{M}}(\mathbf{r})$$

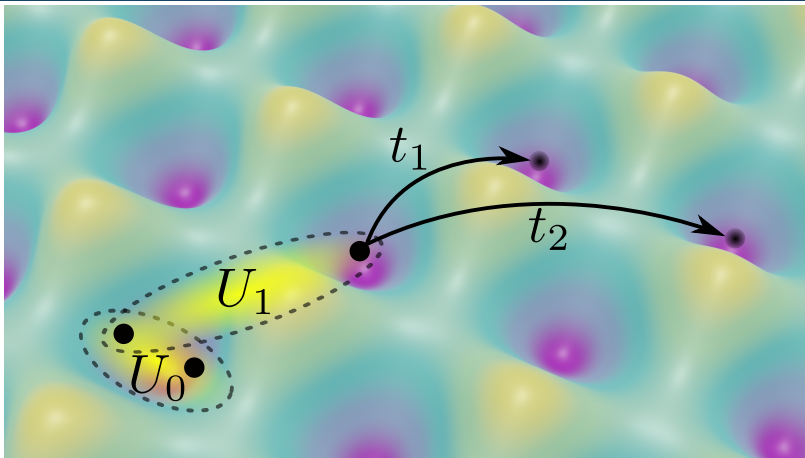
- approximate excitons as bosons
- effective single-particle Hamiltonian for center of mass motion
- solve H in Bloch basis
- Wannier functions from Bloch functions via Fourier like transform

IX Wave Functions

$$H = -\frac{\hbar^2}{2M}\Delta + V^M(\mathbf{r})$$



Bose-Hubbard Model



Götting *et al.*, Phys. Rev. B **105**, 165419 (2022)

Bose-Hubbard Model

$$H = -t \sum_{\langle ij \rangle} \hat{b}_i^\dagger \hat{b}_j + \frac{U}{2} \sum_i \hat{n}_i (\hat{n}_i - 1) - \mu \sum_i \hat{n}_i + \frac{V}{2} \sum_{\langle ij \rangle} \hat{n}_i \hat{n}_j$$

Bose-Hubbard Model

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Mean-field approximation:

$$\varphi_i = \langle \hat{b}_i \rangle$$

$$\varrho_i = \langle \hat{n}_i \rangle$$

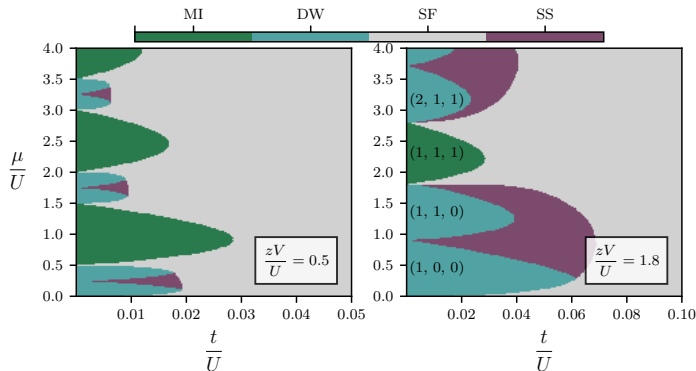
Bose-Hubbard Model

$$H = -t \sum_{\langle ij \rangle} \hat{b}_i^\dagger \hat{b}_j + \frac{U}{2} \sum_i \hat{n}_i (\hat{n}_i - 1) - \mu \sum_i \hat{n}_i + \frac{V}{2} \sum_{\langle ij \rangle} \hat{n}_i \hat{n}_j$$

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Hubbard Parameters

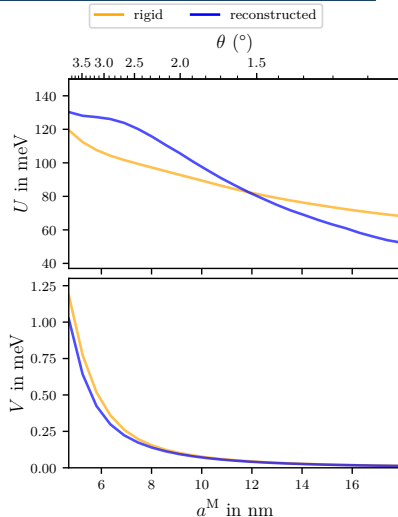
Hopping parameter:

$$t_n = \frac{1}{N} \sum_{\mathbf{Q}} e^{i(\mathbf{R}_0 - \mathbf{R}_n) \cdot \mathbf{Q}} E_{\mathbf{Q}}$$

Hubbard U parameter:

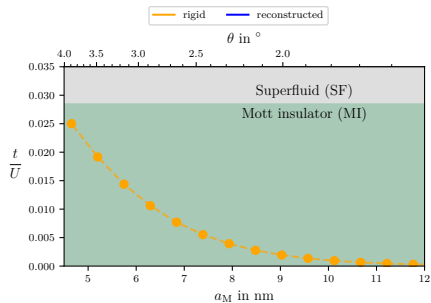
$$U_n = \int \int_{\mathbb{R}^2} |\omega_{\mathbf{R}_0}(\mathbf{r})|^2 |\omega_{\mathbf{R}_n}(\mathbf{r}')|^2 \tilde{U}(\mathbf{r} - \mathbf{r}') d^2r d^2r'$$

hBN is used as sub- and superstrate

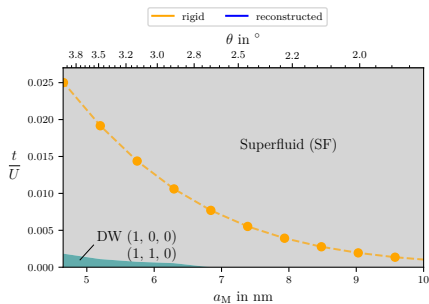


Quantum Phases of IXs

- maximal one boson per lattice site
- material system: MoS₂/WS₂
- relaxed: phases are pushed deeper into MI region



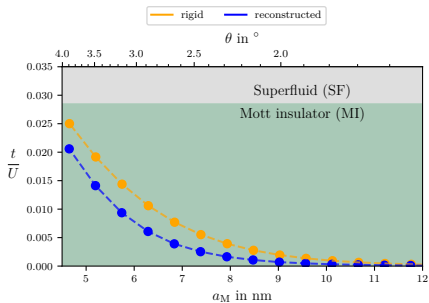
filling: $\rho = 1$



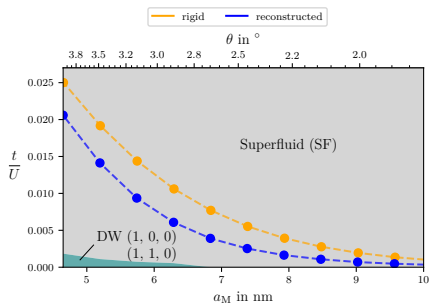
filling: $\rho = \frac{1}{3}, \frac{2}{3}$

Quantum Phases of IXs

- maximal one boson per lattice site
- material system: MoS₂/WS₂
- relaxed: phases are pushed deeper into MI region

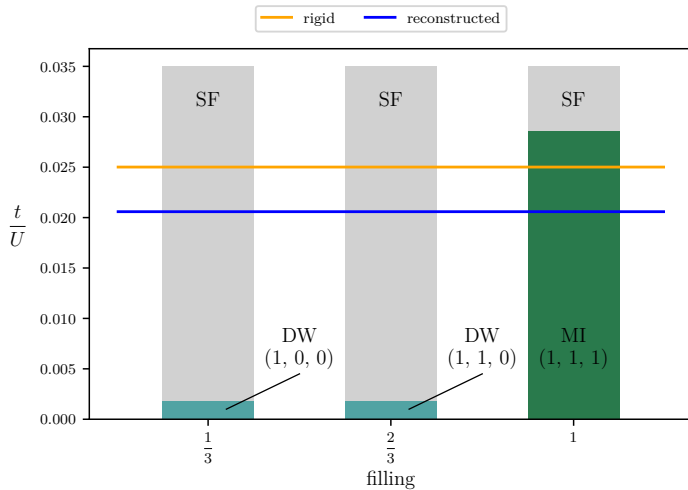


filling: $\rho = 1$

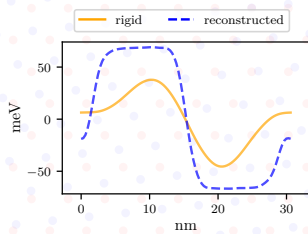


filling: $\rho = \frac{1}{3}, \frac{2}{3}$

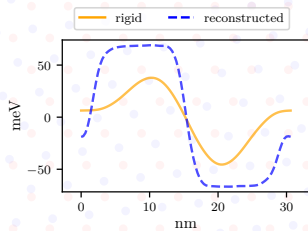
Quantum Phases of IXs



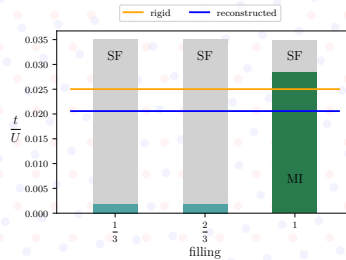
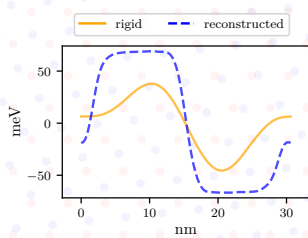
Conclusion



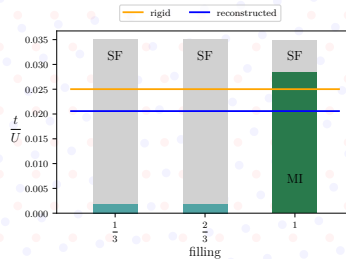
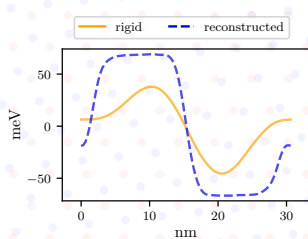
Conclusion



Conclusion



Conclusion



Thank you for listening!

Outlook

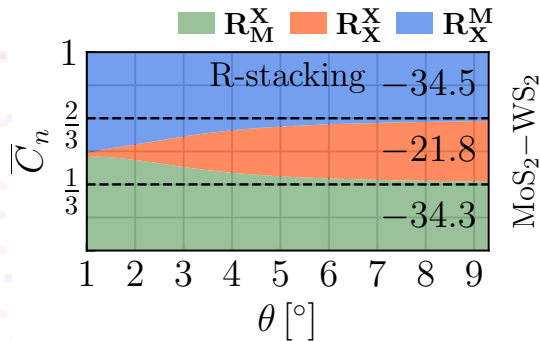
- cluster mean-field [1] or Monte-Carlo methods [2]
- finite-temperature theory by using the thermal average [1]
- include description of exciton-polaritons [3]

[1] Malakar *et al.* Phys. Rev. B 102, 184515 (2020)

[2] Bogner *et al.* Eur. Phys. J. B 92, 111 (2019)

[3] Byrnes *et al.* Phys. Rev. B 81, 205312 (2010)

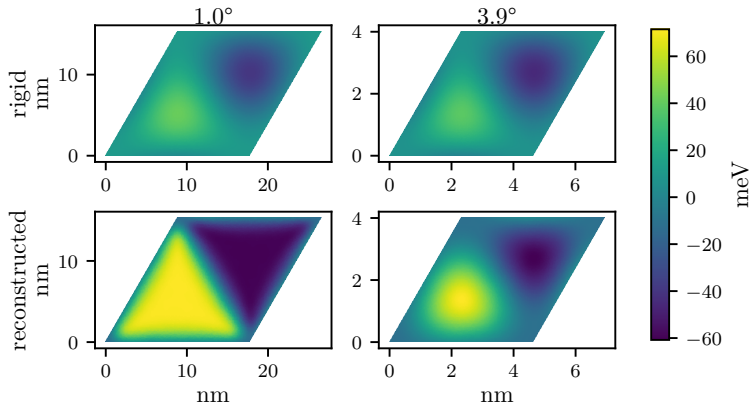
Atomic Reconstructions



Band gap as a function of in-plane distance between adjacent metal atoms and interlayer distance:

$$E_g = E_g(\Delta r_m, d(\Delta r_m))$$

Moiré Potential



IX Hamiltonian

Effective single particle Hamiltonian for IXs:

$$H = -\frac{\hbar^2}{2M}\Delta_{\mathbf{r}} + V^{\text{M}}(\mathbf{r})$$

Bloch functions:

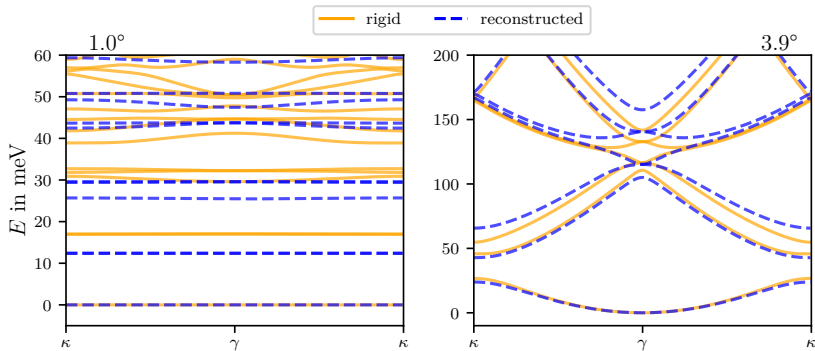
$$\chi_{\mathbf{Q}}^{(\alpha)} = \frac{1}{\sqrt{V}} \sum_{\mathbf{G}^{\text{M}}} c_{\mathbf{Q}-\mathbf{G}^{\text{M}}}^{(\alpha)} e^{i(\mathbf{Q}-\mathbf{G}^{\text{M}})\mathbf{r}}$$

Wannier functions as IX wave functions:

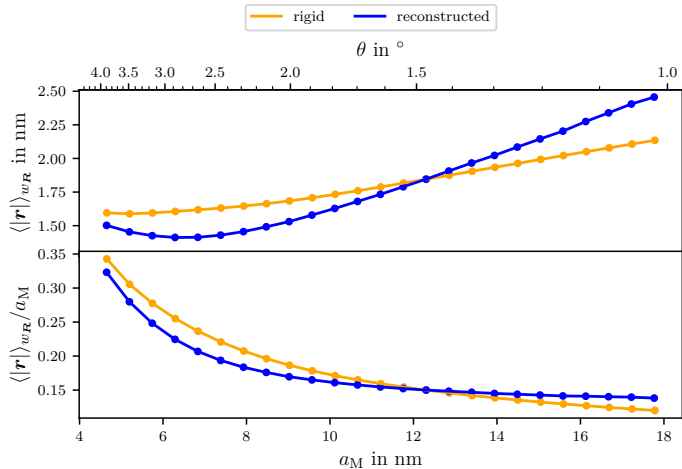
$$\omega_{\mathbf{R}}(\mathbf{r}) = \frac{1}{\sqrt{N}} \sum_{\mathbf{Q}} e^{-i\mathbf{Q}\mathbf{R}} \chi_{\mathbf{Q}}(\mathbf{r})$$

Band Structure

$$\hat{H} = -\frac{\hbar^2}{2M}\Delta + V^M(\mathbf{r})$$



Quantum Phases of IXs



Bose-Hubbard Model

Extended Bose-Hubbard Hamiltonian:

$$H = -t \sum_{\langle ij \rangle} \hat{b}_i^\dagger \hat{b}_j + \frac{U}{2} \sum_i \hat{n}_i (\hat{n}_i - 1) - \mu \sum_i \hat{n}_i + \frac{V}{2} \sum_{\langle ij \rangle} \hat{n}_i \hat{n}_j$$

Mean-field approximation:

$$\varphi_i = \langle \Psi | \hat{b}_i | \Psi \rangle$$

$$\varrho_i = \langle \Psi | \hat{n}_i | \Psi \rangle$$

$$\begin{aligned} \hat{b}_i^\dagger \hat{b}_j &= (\varphi_i^* + \delta \hat{b}_i^\dagger)(\varphi_j + \delta \hat{b}_j) \\ &= \varphi_i^* \varphi_j + \varphi_i^* \delta \hat{b}_j + \varphi_j \delta \hat{b}_i^\dagger + \delta \hat{b}_i^\dagger \delta \hat{b}_j \end{aligned}$$

Bose-Hubbard Model

Extended Bose-Hubbard Hamiltonian:

$$H_{\text{UC}}^{\text{MF}} = -t \sum_X (-\varphi_X^* \bar{\varphi}_X + \bar{\varphi}_X^* \hat{b}_X + \bar{\varphi}_X \hat{b}_X^\dagger) + \frac{U}{2} \sum_X \hat{n}_X (\hat{n}_X - 1) - \mu \sum_X \hat{n}_X \\ + V \sum_X \bar{\varrho}_X (\hat{n}_X - \frac{1}{2} \varrho_X)$$

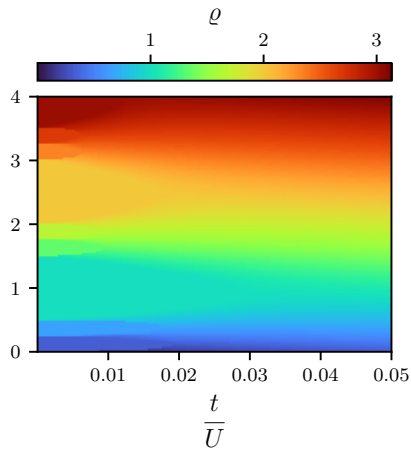
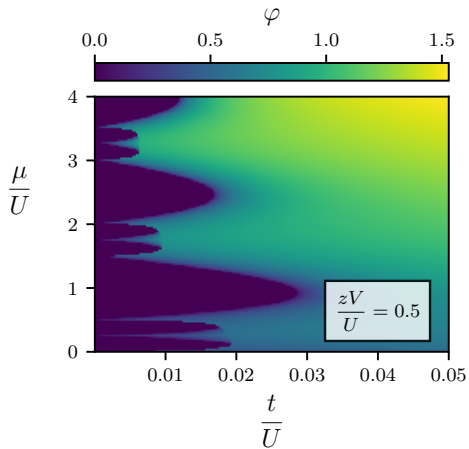
Hopping parameters:

$$t_n = \frac{1}{N} \sum_Q e^{i(\mathbf{R}-\mathbf{R}')\mathbf{Q}} E_Q$$

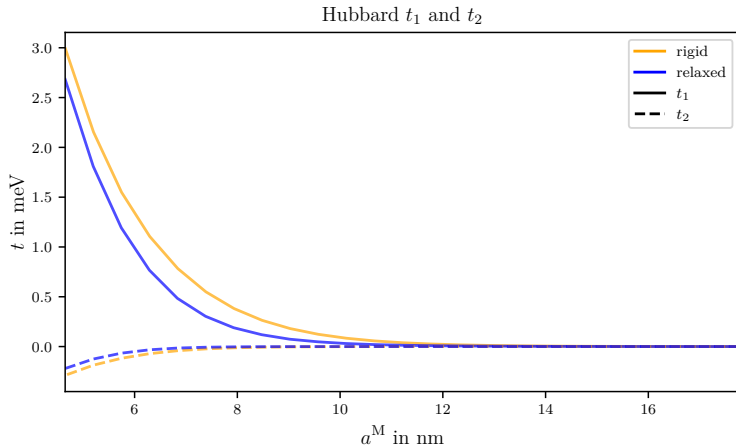
Interaction parameters:

$$U_n = \int \int_{\mathbb{R}^2} |\omega_{\mathbf{R}_0}(\mathbf{r})|^2 |\omega_{\mathbf{R}_n}(\mathbf{r}')|^2 \tilde{U}(\mathbf{r} - \mathbf{r}') d^2r d^2r'$$

Bose-Hubbard Model



Hopping Parameter



Dipole-Dipole Interaction

Dipole-dipole repulsive interaction:

$$\tilde{U}(r) = \frac{e^2}{2\pi\epsilon_r\epsilon_0} \left(\frac{1}{r} - \frac{1}{\sqrt{r^2 + d^2}} \right)$$

Transformation to momentum space:

$$\tilde{U}(q) = \frac{e^2}{\epsilon_0\epsilon_r} \frac{1}{q} \left(1 - e^{-dq} \right)$$

$$\epsilon_r(q) \rightarrow \epsilon_{\text{TMD}} \frac{1 - \tilde{\epsilon}_1\alpha - \tilde{\epsilon}_2\alpha + \tilde{\epsilon}_1\tilde{\epsilon}_2}{1 + \tilde{\epsilon}_1\alpha + \tilde{\epsilon}_2\alpha + \tilde{\epsilon}_1\tilde{\epsilon}_2}$$

$$\tilde{\epsilon}_1 = \frac{1 - \epsilon_{\text{hBN}}}{1 + \epsilon_{\text{hBN}}}, \quad \tilde{\epsilon}_2 = \frac{\epsilon_{\text{TMD}} - 1}{\epsilon_{\text{TMD}} + 1}, \quad \alpha = e^{-qh_{\text{TMD}}}$$

Dipole-Dipole Interaction

Factor accounting the finite width of the bilayer¹

$$F(q) = \frac{2}{\pi} \arctan \left(\frac{\pi}{qh_{\text{eff}}} \right)$$

$$\tilde{U}(\mathbf{r}) = \frac{1}{2\pi} \frac{e^2}{\varepsilon_0 \varepsilon_{\text{TMD}}} \frac{1}{\varepsilon_{\text{TMD}}} \int_0^\infty F(q) \left(1 - e^{-dq}\right) \frac{1 + \tilde{\varepsilon}_1 \alpha + \tilde{\varepsilon}_2 \alpha + \tilde{\varepsilon}_1 \tilde{\varepsilon}_2}{1 - \tilde{\varepsilon}_1 \alpha - \tilde{\varepsilon}_2 \alpha + \tilde{\varepsilon}_1 \tilde{\varepsilon}_2} J_0(qr) dq$$

¹Rösner *et al.* Phys. Rev. B 92, 085102 (2015)

Literature

- Helium on graphene: Yu *et al.*, Phys. Rev. B 103, 235414 (2021)
- Quantum Simulator:
Georgescu *et al.*, Rev. Mod. Phys. 86, 153 (2014)
Feynman, International Journal of Theoretical Physics volume 21, pages 467-488 (1982)
- Reconstructions: Westen *et al.*, Nat. Nanotechnol. 15, 592–597 (2020)