

Local spectra of adaptive domain decomposition methods

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1 Introduction

For second order elliptic partial differential equations, such as diffusion or elasticity, with arbitrary and high coefficient jumps, the convergence rate of domain decomposition methods with classical coarse spaces typically deteriorates. One remedy is the use of adaptive coarse spaces, which use eigenfunctions computed from local generalized eigenvalue problems to enrich the standard coarse space; see, e.g., [19, 6, 5, 4, 22, 23, 3, 16, 17, 14, 7, 8, 24, 1, 20, 2, 13, 21, 10, 9, 11]. This typically results in a condition number estimate of the form

$$\kappa \leq C \text{tol} \quad \text{or} \quad \kappa \leq C \frac{1}{\text{tol}} \quad (1)$$

of the preconditioned system, where C is independent of the coefficient function and tol is a tolerance for the selection of the eigenfunctions.

Obviously, the robustness of the adaptive domain decomposition methods is therefore closely related to the choice of tol . Whereas for a pessimistic choice, i.e., $\text{tol} \approx 1$, the adaptive coarse space can resort to a direct solver, a very optimistic choice can lead to bad convergence behavior of the method.

In this article, we will compare the spectra of the generalized eigenvalue problems of several adaptive coarse spaces for overlapping as well as nonoverlapping domain decomposition methods. The spectra are of interest because they provide information for choosing an adequate tolerance splitting bad and good eigenmodes as well as about the resulting dimension of the adaptive coarse spaces. Therefore, we will consider certain representative examples of coefficient functions in two dimensions.

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Note that we are not going to discuss other important properties of the adaptive coarse spaces considered here, such as

- condition number and iteration counts of the methods,
- costs for the computation of the eigenvalue problems and the coarse basis functions, respectively,
- necessary communication in a parallel implementation and the ratio of local and global work.

Thus, we do not claim to draw a general comparison of the different adaptive methods. We only want to discuss reasonable choices for the user-defined tolerance for different, exemplary coefficient distributions and the different types of eigenvalue problems. We hope that this gives some insight for further discussions.

Model problems and domain decomposition notation We consider the variational form of a second order elliptic partial differential equation, such as diffusion or elasticity, and denote the coefficient by $\rho \in \mathbb{R}^+$ which is assumed to be constant on each finite element. In matrix form, the problem reads $Ax = b$.

Now, let Ω be decomposed into nonoverlapping subdomains $\Omega_1, \dots, \Omega_N$ and Γ be the interface of this domain decomposition. We define corresponding subdomain stiffness matrices $A^{(i)}$ with Neumann boundary conditions on $\partial\Omega_i$, $i = 1, \dots, N$ and the block diagonal matrix $A_N := \text{blockdiag}_i(A^{(i)})$ which is not assembled in the interface degrees of freedom. For an edge \mathcal{E} or its closure $\bar{\mathcal{E}}$ shared by the subdomains Ω_i and Ω_j , we obtain the matrix $A_a^{(i,j)}$ by assembly of the degrees of freedom on \mathcal{E} or $\bar{\mathcal{E}}$, respectively, in the matrix $A_{na}^{(i,j)} := \text{blockdiag}(A^{(i)}, A^{(j)})$.

The Schur complements with respect to $\mathcal{Z} = \mathcal{E}$, $\mathcal{Z} = \bar{\mathcal{E}}$, or any other $\mathcal{Z} \subset \Gamma$ are obtained from $A_{na}^{(i,j)}$ or $A_a^{(i,j)}$ by elimination of all remaining local degrees of freedom \mathcal{Z}^C :

$$S_{*,\mathcal{Z}}^{(i,j)} := A_{*,\mathcal{Z}\mathcal{Z}}^{(i,j)} - A_{*,\mathcal{Z}\mathcal{Z}^C}^{(i,j)} (A_{*,\mathcal{Z}^C\mathcal{Z}^C}^{(i,j)})^{-1} A_{*,\mathcal{Z}^C}^{(i,j)}$$

with $*$ $\in \{a, na\}$. We also need $S_{\mathcal{Z}}^{(i)} := A_{\mathcal{Z}\mathcal{Z}}^{(i)} - A_{\mathcal{Z}\mathcal{Z}^C}^{(i)} (A_{\mathcal{Z}^C\mathcal{Z}^C}^{(i)})^{-1} A_{\mathcal{Z}^C\mathcal{Z}}^{(i)}$.

In addition to that, let the matrices $A_{\mathcal{E}}$ and $M_{\mathcal{E}}$ be matrix discretizations of the one-dimensional bilinear forms $a_{\mathcal{E}}(u, v) := \int_{\mathcal{E}} \rho_{\mathcal{E}, \max} D_{x^t} u D_{x^t} v dx$ and $b_{\mathcal{E}}(u, v) := h^{-1} \sum_{x_k \in \mathcal{E}} \beta_k u(x_k) v(x_k)$ with $\beta_k := \sum_{\{t \in \tau_h: k \in \text{dof}(t)\}} \rho_t$. Here, ρ_t is the constant coefficient on the element $t \in \tau_h$ and $\rho_{\mathcal{E}, \max}(x) := \max \left\{ \lim_{y_i \in \Omega_i \rightarrow x} \rho(y_i), \lim_{y_j \in \Omega_j \rightarrow x} \rho(y_j) \right\}$.

D_{x^t} denotes the tangent derivative with respect to the edge e_{ij} , and the x_k correspond to the finite element nodes on the edge. Consequently, $A_{\mathcal{E}}$ and $M_{\mathcal{E}}$ are the stiffness matrix and a scaled lumped mass matrix on the edge \mathcal{E} .

2 Various adaptive coarse spaces in domain decomposition

Overlapping Schwarz methods We extend the nonoverlapping subdomains to overlapping subdomains $\Omega'_1, \dots, \Omega'_N$ and consider two-level overlapping Schwarz methods of the form

$$M_{OS-2}^{-1} = \Phi A_0^{-1} \Phi^T + \sum_{i=1}^N R_i^T A_i^{-1} R_i,$$

with overlapping matrices $A_i = R_i A R_i^T$, $i = 1, \dots, N$, where R_i is the restriction matrix to the overlapping subdomain Ω'_i , and the coarse matrix $A_0 = \Phi^T A \Phi$. Here, the columns of Φ are the coarse basis functions. We consider three different adaptive coarse spaces for overlapping Schwarz methods, i.e., the Spectral Harmonically Enriched Multiscale (SHEM) [7], the Overlapping Schwarz Approximate Component Mode Synthesis (OS-ACMS) [10], and the Adaptive Generalized Dryja-Smith-Widlund (AGDSW) [9, 11] coarse spaces.

In all these approaches, the coarse space consists of vertex- and edge-based energy-minimizing basis functions, i.e., the interior values Φ_I are given by $\Phi_I := -A_{II}^{-1} A_{I\Gamma} \Phi_\Gamma$ for given interface values Φ_Γ . The vertex-based basis functions are nodal basis functions of Multiscale Finite Element Method (MsFEM) [12] type with different choices of edge values; cf. [7, 10, 9, 11]. The edge-based basis functions are energy-minimizing extensions of the solutions of generalized eigenvalue problems corresponding to the edges of the nonoverlapping domain decomposition.

For an edge \mathcal{E} of the nonoverlapping domain decomposition, we consider the following edge eigenvalue problems.

(Ov1) SHEM coarse space [7]: find $(\tau_\mathcal{E}, \mu_\mathcal{E}) \in V_0^h(\mathcal{E}) \times \mathbb{R}$ s. t.

$$\theta^T A_\mathcal{E} \tau_\mathcal{E} = \mu_\mathcal{E}^{-1} \theta^T M_\mathcal{E} \tau_\mathcal{E} \quad \forall \theta \in V_0^h(\mathcal{E}).$$

(Ov2) OS-ACMS coarse space [10]: find $(\tau_\mathcal{E}, \mu_\mathcal{E}) \in V_0^h(\mathcal{E}) \times \mathbb{R}$ s. t.

$$\theta^T S_\mathcal{E}^{(i,j)} \tau_\mathcal{E} = \mu_\mathcal{E}^{-1} \theta^T A_{\overline{\mathcal{E}\mathcal{E}}} \tau_\mathcal{E} \quad \forall \theta \in V_0^h(\mathcal{E}).$$

(Ov3) AGDSW coarse space [9, 11]: find $(\tau_\mathcal{E}, \mu_\mathcal{E}) \in V_0^h(\mathcal{E}) \times \mathbb{R}$ s. t.

$$\theta^T S_\mathcal{E}^{(i,j)} \tau_\mathcal{E} = \mu_\mathcal{E}^{-1} \theta^T A_{\mathcal{E}\mathcal{E}} \tau_\mathcal{E} \quad \forall \theta \in V_0^h(\mathcal{E}).$$

Let the reciprocal eigenvalues $\mu_\mathcal{E}$ be ordered nondescendingly. Then, we select eigenpairs with $\mu_\mathcal{E} > \text{tol}$ to obtain a condition number estimate of the form $\kappa(M_{OS2}^{-1} A) \leq C \text{tol}$ that is independent of the coefficient function ρ . Note that we use the reciprocal eigenvalue only for comparison with the adaptive coarse spaces for nonoverlapping domain decomposition methods. For the AGDSW coarse space, the matrix on the left hand side is singular. Therefore, we obtain infinity reciprocal eigenvalues in our numerical results.

Nonoverlapping methods In the nonoverlapping domain decomposition methods FETI-1 and FETI-DP, we use the block diagonal matrix A_N and introduce a jump operator B for the interface with $B := (B_1, \dots, B_N)$, $u = (u_1^T, \dots, u_N^T)^T$, and $u_i : \Omega_i \rightarrow \mathbb{R}$, $i = 1, \dots, N$ such that $Bu = 0$ if and only if u is continuous across the interface. The FETI master system is given by

$$\begin{bmatrix} A_N & B^T \\ B & 0 \end{bmatrix} \begin{bmatrix} u \\ \lambda \end{bmatrix} = \begin{bmatrix} f \\ 0 \end{bmatrix}.$$

In FETI-1, the null space of A_N is handled by a projection P such that we solve the following system reduced to the Lagrange multipliers and preconditioned by the nonadaptive, projected Dirichlet preconditioner $PM_D^{-1}P^T$

$$PM_D^{-1}P^T BA_N^+ B^T P^T \lambda = PM_D^{-1}P^T d$$

with corresponding right hand side $PM_D^{-1}P^T d$. We have $M_D^{-1} = B_D \text{blockdiag}(S_{\Gamma_i}^{(i)}) B_D^T$, where B_D is a scaled variant of B . In FETI-DP, we subassemble A_N in a selected number of degrees of freedom on the interface, e.g., all vertices, and denote the resulting nonsingular matrix by \tilde{A}_N . In the nonadaptive case, we then solve the preconditioned system

$$M_D^{-1} B \tilde{A}_N^{-1} B^T \lambda = M_D^{-1} \tilde{d}.$$

Adaptive constraints can then be enforced by, e.g., a second projection P_0 (see [23] for FETI-1 or [17, 14] for FETI-DP) or via a generalized transformation-of-basis approach; see [15]. In FETI-1/-DP and BDD(C) methods, the operator $P_D = B_D^T B$ is used for proving condition number bounds and thus also appears in some generalized eigenvalue problems.

In this paper, we consider the the GenEO eigenvalue problems for FETI-1 (or BDD methods); see [23]; which were first introduced for overlapping Schwarz methods; see [22]. A P_D -based estimate based coarse space was motivated in [19]. There, P_D was localized to $P_{D,\mathcal{E}}$ by extracting from B and B_D the rows only considering the jumps on the corresponding edge (in 2D). A condition number bound for the 2D case was proven in [17]. The method was extended to a robust three dimensional version in [14]. We present results with ρ -scaling as (NOv2a) and deluxe-scaling as (NOv2b). Another P_D -based coarse space was proposed by [3] for BDDC with deluxe-scaling. In the eigenvalue problems, the matrix operator $A : B = A(A+B)^+ B$ is used and the cutoff of the interface Schur complement at the edge $S_{\Gamma_i|\mathcal{E}}^{(i)}$ is used on the right hand side. The energy comparison was generalized to arbitrary scaling matrices $D^{(i)}$ in [17]. Extensions of this method to three dimensions were considered, e.g., in [24, 1, 20, 2, 13]. We present results for ρ -scaling as (NOv3a) and deluxe-scaling as (NOv3b).

(NOv1) GenEO coarse space (FETI-1/BDD) [23]: find $(\tau_{\Gamma_i}, \mu_{\Gamma_i}) \in V^h(\Gamma_i) \times \mathbb{R}$ s. t.

$$\theta^T S_{\Gamma_i}^{(i)} \tau_{\Gamma_i} = \mu_{\Gamma_i}^{-1} \theta^T (B_i^T M_D^{-1} B_i) \tau_{\Gamma_i} \quad \forall \theta \in V^h(\Gamma_i).$$

(NOv2) P_D -based coarse space no. 1 (FETI-DP/BDDC) [19]: find $(\tau_{\Gamma_i}, \mu_{\Gamma_i}) \in (\ker S_{na,\Gamma_{ij}}^{(i,j)})^\perp \times \mathbb{R}$ s. t.

$$\theta^T P_{D,\mathcal{E}}^T S_{na,\Gamma_{ij}}^{(i,j)} P_{D,\mathcal{E}} \tau_{\Gamma_{ij}} = \mu_{\Gamma_{ij}} \theta^T S_{na,\Gamma_{ij}}^{(i,j)} \tau_{\Gamma_{ij}} \quad \forall \theta \in (\ker S_{na,\Gamma_{ij}}^{(i,j)})^\perp.$$

(NOv3) P_D -based coarse space no. 2 (FETI-DP/BDDC) [3]: find $(\tau_{\mathcal{E}}, \mu_{\mathcal{E}}) \in V_0^h(\mathcal{E}) \times \mathbb{R}$ s. t.

$$\theta^T S_{\mathcal{E}}^{(i)} : S_{\mathcal{E}}^{(j)} \tau_{\mathcal{E}} = \mu_{\mathcal{E}} \theta^T (D_{\mathcal{E}}^{(j),T} S_{\Gamma_i|\mathcal{E}}^{(i)} D_{\mathcal{E}}^{(j)} + D_{\mathcal{E}}^{(i),T} S_{\Gamma_i|\mathcal{E}}^{(j)} D_{\mathcal{E}}^{(i)}) \tau_{\mathcal{E}} \quad \forall \theta \in V_0^h(\mathcal{E})$$

Let the (reciprocal) eigenvalues be ordered nondecreasingly. Then, we select eigenpairs with $\mu_{\Gamma_i}^{-1}$, $\mu_{\Gamma_{ij}}$, or $\mu_{\mathcal{E}}$ greater than tol . For the (NOv1) and the (NOv3) eigenvalue problems, the matrix on the left hand side is singular, therefore, we obtain infinity (reciprocal) eigenvalues in our numerical results. For (NOv1), note that the authors of [23] do not incorporate the eigenvectors corresponding to zero eigenvalues into the coarse space. With all three eigenvalue problems (NOv1)-(NOv3), we then obtain adaptive methods with a condition number bound $\kappa \leq C\text{tol}$ that is independent of the coefficient function ρ .

3 Numerical results

In this section, we present results for a diffusion problem on $\Omega = (0, 1)^2$ decomposed into nine subdomains. We used a rectangular domain decomposition and slightly curved edges for the subdomain in the center to prevent the appearance of symmetric effects. We set homogeneous Dirichlet boundary conditions for the edge with $x = 0$ and homogeneous Neumann boundary conditions elsewhere.

The local spectra of the different adaptive coarse spaces for eight different coefficient distributions are shown in Figures 1 and 2. The critical eigenvalues and reciprocal eigenvalues, respectively, are displayed above the spectral gap, which is hatched in gray. They are plotted side by side if they are close to each other. A wide spectral gap simplifies the choice of an appropriate tolerance tol . In addition to that, the number of critical eigenvalues is related to the dimension of the coarse space. Note that the condition number estimate (1) guarantees fast convergence of all different approaches for arbitrary coefficient distributions if a suitable tolerance is chosen. However, as can be observed from our results, there are significant differences in the width of the spectral gap and the number of critical eigenvalues for the depicted model problems.

The use of harmonic extensions in the eigenvalue problems of the OS-ACMS coarse space can reduce the number of bad eigenmodes compared to the cheaper one-dimensional integrals in the related SHEM coarse space. A similar behavior can be observed for the expensive deluxe-scaling compared to the cheaper ρ -scaling for the P_D -based approaches for FETI-DP/BDDC. For several coefficient distributions, the width of the spectral gap is larger than two orders of magnitude for all approaches, whereas it is quite small, e.g., for channel-type coefficient distributions.

Note that the plots in Figures 1 and 2 contain much more information, which we cannot discuss here due to lack of space. We hope that the results presented here give some insight for further investigations. Further investigations in three dimensions are also of high interest. This is however out of the scope of this paper. A small comparison between the 3D version algorithms of columns (Nov2a) and (Nov2b) can be found in [18, Sec. 6.5.3]. For overlapping Schwarz methods, a comparison between different 3D approaches including (Ov3) can be found in [11].

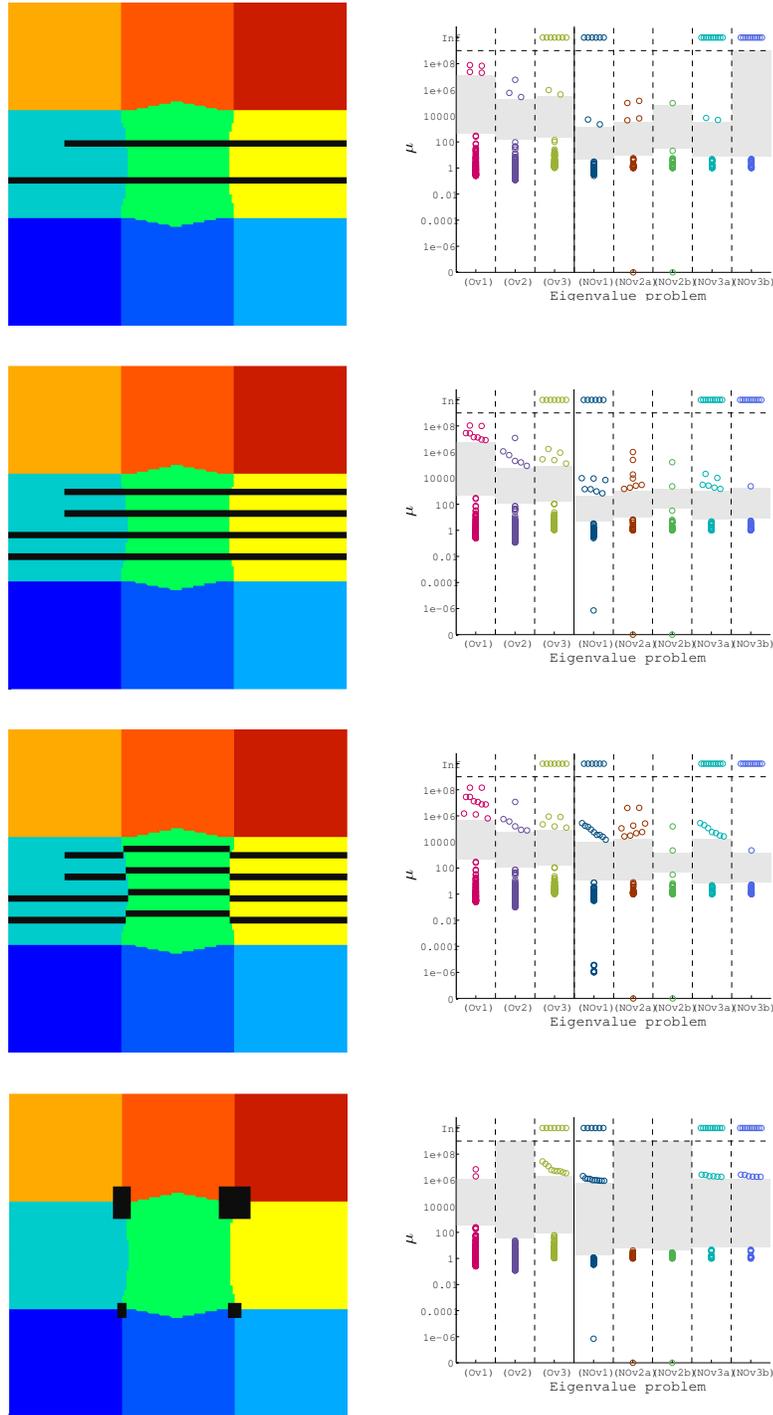


Fig. 1 For each exemplary coefficient distribution, domain decomposition and coefficient distribution is shown on the left hand side. Large coefficients with $\rho = 1e6$ are shown in black (low coefficients with $\rho = 1$ are not shown). The different subdomains are shown in different colors on a layer underneath the large coefficients. The corresponding (reciprocal) eigenvalues μ are shown on the right hand side. Large values (greater 500) are distributed horizontally within the columns to visualize their number. The gap between good and bad eigenmodes is shown in gray.

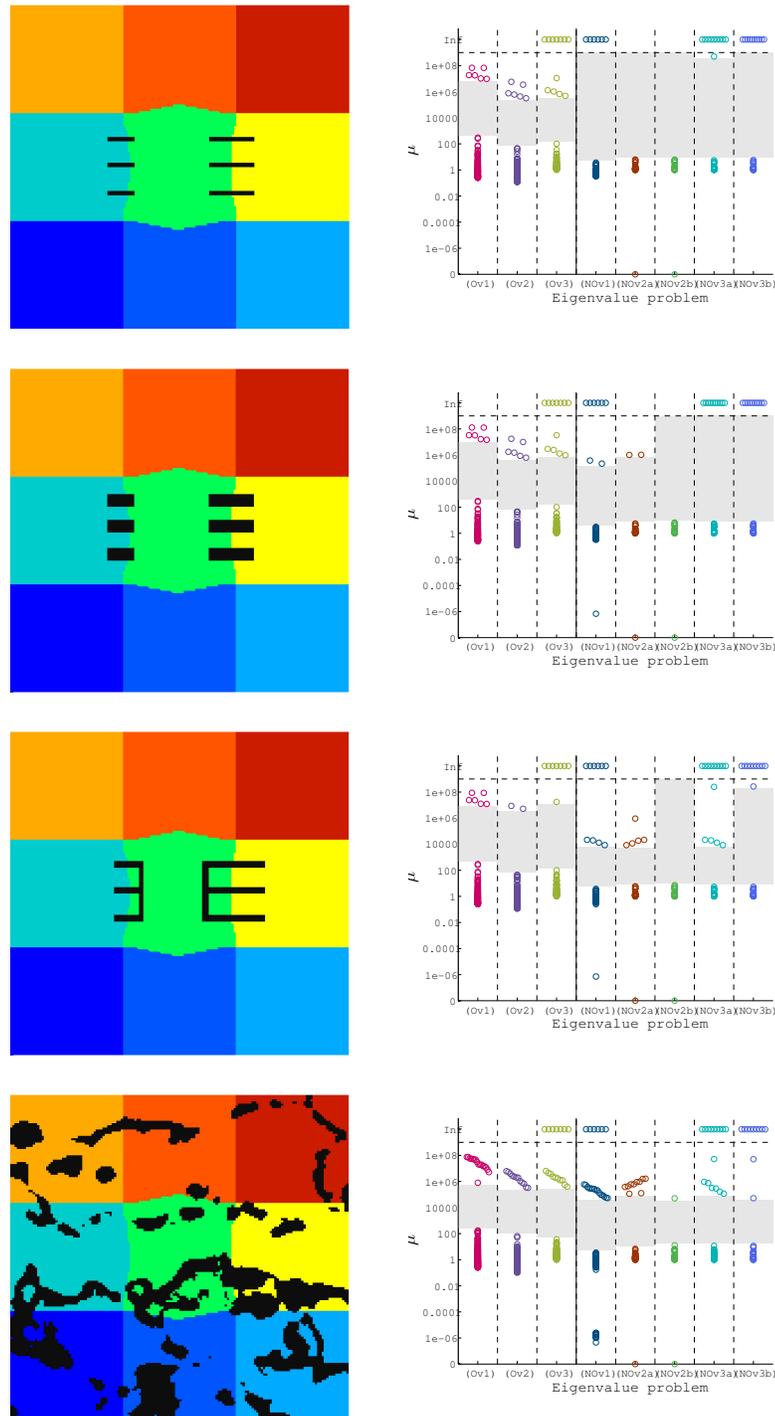


Fig. 2 For each exemplary coefficient distribution, domain decomposition and coefficient distribution is shown on the left hand side. Bottom image: coefficient function generated from the microstructure of a dual-phase steel; courtesy of Jörg Schröder, University of Duisburg-Essen, Germany, originating from a cooperation with ThyssenKruppSteel. Large coefficients with $\rho = 1e6$ are shown in black (low coefficients with $\rho = 1$ are not shown). The different subdomains are shown in different colors on a layer underneath the large coefficients. The corresponding (reciprocal) eigenvalues μ are shown on the right hand side. Large values (greater 500) are distributed horizontally within the columns to visualize their number. The gap between good and bad eigenmodes is shown in gray.

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