Long-term monitoring of Ark 120 with Swift

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ABSTRACT

We report the results of a six-month Swift monitoring campaign of Ark 120, a prototypical 'bare' Seyfert 1 galaxy. The lack of intrinsic absorption combined with the nearly contemporaneous coverage of the ultraviolet (UV) and X-ray bands makes it possible to investigate the link between the accretion disc and the putative Comptonization corona. Our observations confirm the presence of substantial temporal variability, with the X-ray characterized by large-amplitude flux changes on time-scales of few days, while the variations in the UV bands are smoother and occur on time-scales of several weeks. The source also shows spectral variability with the X-ray spectrum steepening when the source is brighter. We do not detect any correlation between the UV flux and the X-ray spectral slope. A cross-correlation analysis suggests positive delays between X-rays and the UV emission, favouring a scenario of disc reprocessing. Although the strength of the correlation is moderate with a delay which is not well constrained (7.5 \pm 7 d), it is nevertheless indicative of a very large disc reprocessing region, with a separation between the X-ray and the UV-emitting regions, which could be as large as 1000 r_G . The Ark 120 correlation results are in agreement with those obtained in similar multiwavelength monitoring studies of active galactic nuclei (AGN). When combined together, the observations so far can be well described by a linear relation between the X-ray/UV delays and the mass of the central black hole. Within the context of the simplest scenario, where these delays correspond to light-travel times, the implied distance between the X-ray source and the optical/UV disc reprocessing region in these AGN should be of the order of many hundreds of gravitational radii.

Key words: galaxies: active – galaxies: nuclei – X-rays: galaxies.

1 INTRODUCTION

Numerous observational studies across the electromagnetic spectrum have shown that active galactic nuclei (AGN) are powerful emitters of variable radiation from the radio band to γ -ray energies and that a sizable amount of this radiation is emitted in the optical, ultraviolet (UV), X-ray energy bands. It is generally thought that AGN are powered by accretion on to a central supermassive black hole (BH) and that the optical–UV emission is the thermal radiation produced directly from the accretion flow, whereas the X-rays are produced through the Comptonization process in a putative corona. While this general picture is widely accepted, the details of the interaction between disc and corona, the geometry and physical state

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of the latter, as well as the origin of the variability are still poorly understood.

One of the most promising approaches to shed some light on the inner region of BH systems is to track the variable behaviour of the different components of the central engine using multiwavelength monitoring observations. Over the years, coordinated optical and X-ray monitoring observations of AGN have revealed that, in addition to the ubiquitous X-ray variability, long-term optical variations are consistently detected. The latter are generally interpreted as intrinsic variations of the disc emission caused by disturbances propagating across the disc or, alternatively, as a result of the disc reprocessing of the X-ray emission produced by the corona. In this context, various monitoring studies have yielded contrasting results, with some sources showing highly correlated optical/X-ray behaviour with one band leading the other, and some others showing no correlation at all (e.g. Maoz, Edelson & Nandra 2000; Uttley et al. 2003; Arévalo et al. 2008, 2009; Breedt et al. 2009, 2010).

Some of the inferred discrepancies may be attributed to the intrinsic difficulties of coordinating satellite-based X-ray monitoring with ground-based optical observations, which may also be hampered by weather conditions.

An important step forward in this field has been the advent of Swift, which, thanks to its flexibility and simultaneous coverage of several bands in the optical/UV range and in the X-rays, eliminates all possible ambiguities associated with weather conditions and lack of coordination in the different energy bands. In principle, true contemporaneous observations in several energy bands should make it possible to test whether the diverse temporal behaviour observed in different objects is related to some fundamental properties of the central engine. More specifically, different combinations of BH mass $M_{\rm BH}$ and accretion rate \dot{m} may be responsible for the diverse behaviour observed in various BH systems. For example, Uttley et al. (2003) hypothesized that objects with large $M_{\rm BH}$ and low \dot{m} should yield a relatively tight X-ray/optical correlation, since they have cooler discs and hence the optical region closer (in units of gravitational radii) to the X-ray-emitting central corona. Conversely, for AGN with small $M_{\rm BH}$ and high \dot{m} , the region producing the optical is relatively distant from the X-ray region (the corona) and hence an uncorrelated optical/X-ray behaviour may be expected.

In our previous work, we investigated with *Swift* the long-term behaviour of PKS 0558–504, an X-ray bright radio-loud Narrowline Seyfert 1 with $M_{\rm BH} \sim 3 \times 10^8 \,\rm M_{\odot}$ likely accreting at super-Eddington rate. The main findings of our 1.5-yr monitoring campaign can be summarized as follows. PKS 0558–504 is highly variable at all wavelengths probed by *Swift*, with variability levels decreasing from the X-rays to the optical bands. The large-amplitude variations measured by the *UVW2* filter are strongly correlated with all the other optical and UV bands and weakly (but significantly) correlated with the X-ray variations. These results, combined with suggestive evidence, which in PKS 0558–504 perturbations propagate from the outer to the inner parts of the accretion flow and to the corona, confirm the existence of physical link between disc and corona and disfavour the reprocessing scenario (for more details see Gliozzi et al. 2013).

Here, to further investigate the link between accretion disc and corona and its putative dependence on $M_{\rm BH}$ and \dot{m} , we study the correlated variability properties of the prototypical bare Seyfert 1 galaxy Ark 120, which has a relatively large BH mass, $M_{\rm BH} \sim 1.5 \times 10^8 {\rm M}_{\odot}$ (Peterson et al. 2004) and low accretion rate $L_{\rm bol}/L_{\rm Edd} \sim 0.05$ (Vasudevan & Fabian 2007). Past studies have shown that Ark 120 stands out among bright Seyfert galaxies for the lack of any evidence for reddening in infrared and UV observations (Ward, et al. 1987; Crenshaw et al. 1999) and warm absorbers (Vaughan et al. 2004), and is a highly variable source over the entire spectrum (e.g. Kaspi et al. 2000; Carini, Noble & Miller 2003; Doroshenko, Sergeev & Pronik 2008), making it an ideal target for a *Swift* monitoring campaign.

This paper is structured as follows. In Section 2, we describe the observations and data reduction. In Section 3, we study the temporal and spectral variability properties with model-independent tools; in Section 4, we describe the results of a standard spectral analysis, whereas in Section 5 we carry out an inter-band correlation analysis. In Section 6, we summarize our main findings and discuss their implications.

Hereafter, we adopt a cosmology with $H_0 = 71 \text{ km s}^{-1} \text{ Mpc}^{-1}$, $\Omega_{\Lambda} = 0.73$, and $\Omega_{M} = 0.27$ (Bennet et al. 2003). With the assumed cosmological parameters, the luminosity distance of Ark 120 is 142 Mpc.

2 OBSERVATIONS AND DATA REDUCTION

Ark 120 was observed by Swift (Gehrels 2004) between 2014 September 4 and 2015 March 9 with a cadence of one pointing every two days. This observing strategy was devised to preserve the UV/Optical Telescope (UVOT) filter wheel by using the filters of the day U and UVM2, which regularly alternate every two days. The filters choice is based on the fact that U is the 'reddest' among standard filters and UVM2 the 'bluest' clean filter (UVW2 is slightly bluer than UVM2 with a central wavelength of 1950 Å versus 2200 Å, but is affected by a red leak). The details of the Swift monitoring campaign are summarized in Tables 1 (dates and exposures) and 2 (X-ray count rates and hardness ratios, as well as U and UVM2 fluxes), where only the first four entries are shown for illustrative purposes. The complete tables are available in electronic format. The Swift X-Ray Telescope (XRT; Burrows et al. 2005) observations were performed in windowed timing mode to avoid possible pile-up effects (Hill et al. 2004). For both spectra and light curves, source photons were extracted in a circular region with a radius of 20 pixels (corresponding to \sim 47 arcsec) centred on the source, and the background was selected using a circle of the same size but shifted along the window, away from the source position. Only single to quadruple events in the energy range of 0.3-10 keV were selected for further analysis. Source and background spectra were extracted from the event file by using XSELECT version 2.3. All the light curves analysed in Section 3 are background subtracted. The auxiliary response files were created by the SWIFT tool xrtmkarf and used in combination with the response matrix swxwt0to2s6psf1_20131212v001.rmf. We also obtained photometry with the UVOT (Poole et al. 2008; Breeveld et al. 2010) in the U and UVM2. Source photons were extracted from a circular region with r = 5 arcsec, and the background from a sourcefree circular region with a radius of 20 arcsec. The UVOT tool uvotsource was used to determine the fluxes. The fluxes were corrected for Galactic reddening (E(B-V) = 0.113 obtained from NASA/IPAC Extragalactic Database) with the standard reddening correction curves by Cardelli, Clayton & Mathis (1989). The U and UVM2 images of Ark 120 show point-like sources, suggesting that the galaxy contribution is negligible compared to the AGN emission. Indeed, this conclusion is confirmed by Koss et al. (2011), who studied the properties of the host galaxies of the AGN observed with Swift Burst Alert Telescope (BAT) and concluded that the galaxy emission of Ark 120 is strongly contaminated by the AGN in all the filters.

3 VARIABILITY ANALYSIS OF ARK 120

Fig. 1 shows the UVOT U ($\lambda = 344$ nm; top panel), UVM2($\lambda = 231$ nm; middle panel) light curves, as well as the 0.3–10 keV XRT light curve (bottom panel) of Ark 120. On top of the X-ray time series, we superimposed the normalized UVM2 light curve (blue continuous line) to highlight the similarities and differences between the temporal behaviour in the UV and X-rays. While the overall trend seems to be broadly consistent, the X-ray light curve exhibits short-term variability, which is not present in the smoother UV light curves. The dot–dashed and dashed lines on the X-ray light curve represent the threshold values used for the flux-selected spectral analysis (see Section 4); data points above the dot–dashed line have been selected for the 'high-flux' spectrum, whereas those below the dashed line define the 'low-flux' spectrum.

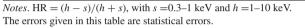
As suggested by the visual inspection of this figure, Ark 120 is highly variable at all wavelengths probed by *Swift*. According to a

Segment	Start time (UT)	End time (UT)	MJD	Observing time given in s		
				$T_{\rm XRT}$	T_U	T_{UVM2}
01	2014-09-04 01:20:00	2014-09-04 01:36:59	56904.56	992	998	_
02	2014-09-06 04:34:56	2014-09-06 04:51:58	56906.69	1009	_	1000
03	2014-09-08 06:03:36	2014-09-08 06:20:58	56908.75	1026	1019	-
04	2014-09-10 10:50:37	2014-09-10 11:07:58	56910.95	1038	_	1018

Table 1.Observation log of Ark 120.

 Table 2. Swift XRT count rates and HR and UVOT fluxes (mJy) of Ark 120.

Segment	XRT rate	XRT HR	U	UVM2
01	1.19 ± 0.03	0.20 ± 0.02	6.93 ± 0.15	_
02	1.25 ± 0.04	0.11 ± 0.02	_	3.57 ± 0.08
03	1.13 ± 0.03	0.15 ± 0.02	6.96 ± 0.15	-
04	1.31 ± 0.04	0.18 ± 0.02	-	3.48 ± 0.08



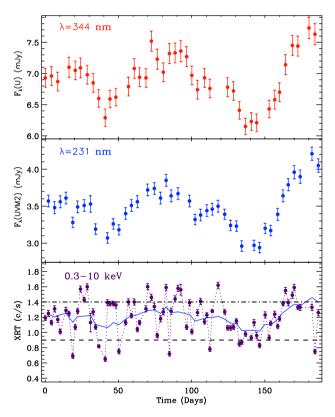


Figure 1. *Swift* UVOT U (top panel), UVM2 (middle panel) and 0.3–10 keV XRT (bottom panel) light curves of Ark 120 from 2014 September 4 to 2015 March 9. The UV fluxes are corrected for Galactic absorption and expressed in units of mJy. On top of the X-ray time series, we superimposed the UVM2 light curve normalized to have the same mean as the X-ray light curve. The dot–dashed and dashed lines on the X-ray light curve represent the threshold values used for the flux-selected spectral analysis.

 χ^2 test, the light curves show significant variability with χ^2 /dof values of 270.5/40, 562.6/41, and 4283.8/80 for the *U*, *UVM2* UVOT filters, and XRT, respectively. An analysis of the fractional variability $F_{\text{var}} = \sqrt{\sigma^2 - \Delta^2} / \langle r \rangle$ (where σ^2 is the variance, Δ^2 the mean square value of the uncertainty associated with each individual count rate, and $\langle r \rangle$ the unweighted mean count rate) confirms these

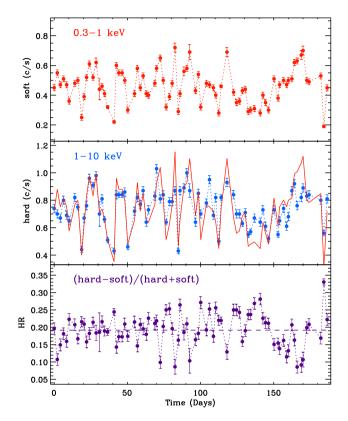


Figure 2. Soft, hard, and hardness ratio, HR = (h - s)/(h + s) light curves Ark 120. The solid line superimposed on the middle panel represents the normalized soft X-ray light curve. The dashed line in the bottom panel describes the average value of the hardness ratio.

results, suggesting the presence of a positive trend between variability and energy band: $F_{\text{var},U} = (5.2 \pm 0.4)$ per cent, $F_{\text{var},UVM2} =$ (7.9 ± 0.4) per cent, and $F_{\text{var, XRT}} = (19.6 \pm 0.3)$ per cent. Fig. 2 shows the soft (0.3-1 keV), hard (1-10 keV), and hardness ratio HR = (h - s)/(h + s) light curves of Ark 120. To guide the eye, we have superimposed the normalized soft X-ray light curve (red continuous line in the middle panel) on top of the hard X-ray light curve. The similar trend between the two X-ray bands indicates that soft and hard time series vary significantly and nearly in concert. According to a χ^2 test and fractional variability analysis, the count rate variation is highly significant with $\chi^2/dof = 2693.5/80$, $F_{var, soft} = (24.9 \pm 0.5)$ per cent, and 4283.8/80, $F_{\text{var, hard}} = (17.9 \pm 0.4)$ per cent for the soft and hard bands, respectively. Based on the same tests, the variability of the HR light curve appears to be statistically significant with χ^2 /dof = 801.3/80 and $F_{\text{var. HR}} = (23 \pm 1)$ per cent, suggesting the presence of X-ray spectral variability.

A simple model-independent method to test for spectral variability makes use of the plot of the X-ray hard count rate versus the soft count rate (e.g. Churazov, Gilfanov & Revnivtsev

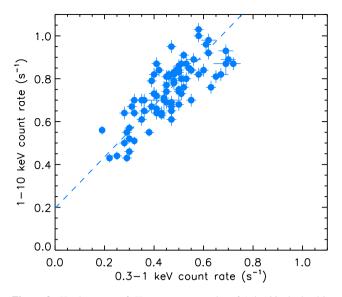


Figure 3. Hard versus soft X-ray count rate plot of Ark 120 obtained in the *Swift* XRT campaign. The dashed line represents the best-fitting linear model CountRate_{1-10 keV} = $0.20 \pm 0.02 + (1.10 \pm 0.04)$ CountRate_{0.3-1 keV}.

2001), which is shown in Fig. 3. As expected from the visual inspection of the X-ray light curves, the plot of hard versus soft count rates shows a strong linear correlation - the Spearman's rank correlation coefficient is $\rho = 0.79$ and associated probability of random correlation $P_{\rho} = 1.4 \times 10^{-18}$ and is well described by the equation CountRate_{1-10keV} = $0.20 \pm 0.02 +$ (1.10 ± 0.04) CountRate_{0.3-1 keV}, obtained with the routine fitexy (Press et al. 1997), which accounts for the errors on both y and xaxes, and will be adopted in the rest of this paper for any linear correlation analysis. Interestingly, while the slope is roughly consistent with the unity, the positive intercept, inconsistent with zero, suggests the existence of a non-variable hard component. An additional direct way to study the X-ray spectral variability of AGN is to plot the hardness ratio versus the total flux. The result of this analysis for Ark 120 is shown in Fig. 4, which, despite the substantial scatter, reveals the existence of a shallow but robust anticorrelation ($\rho = -0.38$, $P_{\rho} = 5 \times 10^{-4}$) described by HR = $0.287 \pm 0.008 - (0.080 \pm 0.007)$ Count Rate_{0.3-10 keV}, which indicates that the spectrum softens as the source brightens.

Finally, model-independent information about the broadband spectral variability can be inferred by studying the temporal evolution of the broad-band spectral index $\alpha_{OX} = \log (l_{2500 \text{ Å}}/l_{2 \text{ keV}})/\log (\nu_{2500 \text{ Å}}/\nu_{2 \text{ keV}})$ (Tananbaum et al. 1979). We derived α_{OX} from the simultaneous X-ray and *UVM2* fluxes, and plotted the light curve in Fig. 5, which suggests the presence of a weak variability of the spectral energy distribution (SED): $\chi^2/\text{dof} = 62.6/39$, $F_{\text{var},\alpha_{OX}} = (1.4 \pm 0.4)$ per cent.

In summary, the six-month *Swift* monitoring campaign of Ark 120 confirms the presence of significant large-amplitude variability (for such a large BH mass) in all bands probed by the UVOT and XRT, with the X-ray band being by far the most variable component, and indicate that the temporal variability of Ark 120 is associated with spectral changes of the X-rays and, to a lesser extent, of the broad-band SED.

4 X-RAY SPECTRAL ANALYSIS

The X-ray spectral analysis was performed using the xSPEC v.12.9.0 software package (Arnaud 1996). For ~ 1 ks exposure observa-

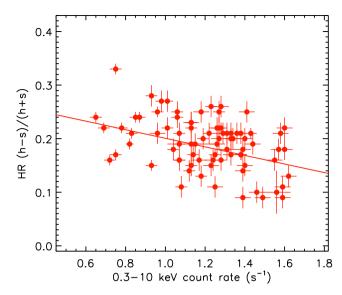


Figure 4. Hardness ratio HR = (h-s)/(h+s) versus 0.3–10 keV X-ray count rate plot of Ark 120 obtained in the *Swift* XRT campaign. The continuous line represents the best-fitting linear model $HR = 0.287 \pm 0.008 - (0.080 \pm 0.007)$ CountRate_{0.3–10 keV}.

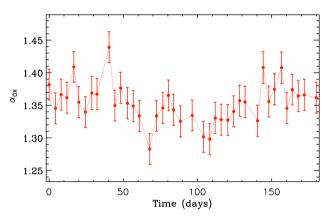


Figure 5. Light curve of the broad-band spectral index α_{OX} .

tions, spectra were rebinned within GRPPHA 3.0.0 to have at least one photon per bin and fitted with the C-statistic, whereas combined flux-selected spectra were rebinned at 20 counts per channel for the χ^2 statistic to be valid. The errors on spectral parameters represent the 68 per cent confidence level (1 σ) for one interesting parameter ($\Delta \chi^2 = 1$). We verified that the two UV data points are well above the extrapolation of the X-ray best-fitting model and are most likely associated with the accretion disc emission. We did not include the UV data in the spectral fitting analysis, because two non-simultaneous data points in the U and UVM2 filters are not sufficient to characterize the properties of the accretion disc.

We carried out a systematic spectral analysis of every observation of the Ark 120 campaign, even though the short exposures of individual *Swift* XRT pointings yield X-ray spectra with limited statistics. We adopted a baseline model which comprises two Comptonization components (representing the primary emission produced by the corona and the soft excess) and a Gaussian line to account for the iron K α line emission. All additive spectral components are absorbed by a column density fixed at the Galactic value $N_{\rm H} = 1.01 \times 10^{21} \,{\rm cm}^{-2}$, parametrized by the wabs model in xspec. This model choice was guided by the spectral results from past studies with higher signal-to-noise data, and more specifically by a recent study based on long exposures from *XMM–Newton* and *NuS-TAR*, which confirmed the presence of a soft excess, which appears to be consistent with an additional cooler Comptonization component (Matt et al. 2014). At the beginning, all parameters are left free to vary. However, given the limited statistics of the spectra, some parameters are poorly constrained and yield unreasonable values when computing their statistical uncertainty; in those cases, the parameters are fixed at their best-fitting value. The parameters left free during the error calculation are the spectral index and the normalization for the individual observations. For the combined spectra of high- and low-flux cases, more parameters are reasonably defined and therefore can be left free to vary during the error calculations.

Both the soft excess and the coronal emission have been parametrized by the Bulk Motion Comptonization (BMC) model in XSPEC (Titarchuk, Mastichiadis & Kylafis 1997), which is a simple but comprehensive Comptonization model which can fit both thermal and bulk Comptonization processes, and is described by four parameters: kT (the temperature of the thermal seed photons), α (the energy spectral index related to the photon index by the relationship $\Gamma = 1 + \alpha$), log (A) (a parameter describing to the Comptonization fraction f = A/(1 + A)), and the normalization. We used the BMC model instead of the phenomenological power-law model because the BMC parameters are computed in a self-consistent way, and the power law produced by BMC does not extend to arbitrarily low energies.

For illustrative purposes of individual XRT spectra yielded during the Ark 120 campaign, the unfolded spectrum (eeufspec in XSPEC) and the data-to-model ratio from obsid 34 with net exposure of 1039 s and count rate of ~ 1.5 c s⁻¹ are shown in Fig. 6. This represents one of the best-case scenarios, since it refers to an observation with relatively long exposure and high count rate. Larger uncertainties are associated with observations with shorter exposures or lower count rates.

All individual observations are reasonably well fitted with this baseline model (χ^2_{red} ranges from 0.6 to 1.14), although only a few spectra statistically require more than one BMC component, suggesting that the model overparametrizes the low signal-to-noise spectra. Not surprisingly, the model parameters are poorly constrained. Nevertheless, since the model-independent analysis suggests the presence of spectral variability throughout the monitoring campaign, we tested whether this finding can be confirmed by constructing a light curve of photon index, describing the primary X-ray emission. The resulting plot, shown in Fig. 7, suggests that the time series of the photon index is consistent with the hypothesis of constancy because of the large uncertainties associated with the Γ values. This is indeed confirmed by a χ^2 test, which yields $\chi^2/dof = 52.6/80 \ (P_{\chi}^2 = 0.99)$. Note that the same conclusion is reached using a single BMC model or a power-law model to fit the continuum.

In an attempt to test whether the spectrum of Ark 120 genuinely steepens when the source brightens, we combined several individual spectra of observations with low count rate (count rate_{0.3 - 10keV} \leq 0.9 c s⁻¹) to produce a 'low-flux' spectrum, and similarly several spectra with high count rate (count rate_{0.3 - 10keV} \geq 1.4 c s⁻¹) to obtain a 'high-flux' spectrum. The threshold count rate values, shown in the bottom panel of Fig. 1, were arbitrarily chosen to be distinct from the mean count rate (1.20 \pm 0.03 c s⁻¹) and to encompass at least 10 individual observations each, in order to increase the signal-to-noise ratio (S/N) of the combined spectra.

The resulting low-flux and a high-flux spectra have well-separated mean count rates, 0.77 \pm 0.02 and 1.54 \pm 0.02 c s^{-1},

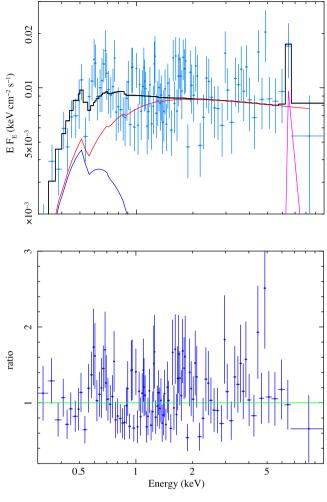


Figure 6. Top panel: unfolded XRT spectrum of Ark 120, obtained using the EEUFSPEC command in XSPEC. The model includes two BMC components plus one Gaussian line modified by photoelectric absorption. Bottom panel: data-to-model ratio.

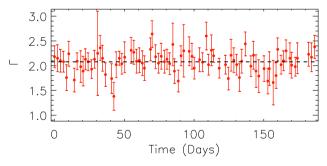


Figure 7. Light curve of the X-ray primary emission photon index Γ during the Ark 120 campaign.

and net exposures of ~11 and ~12 ksec, respectively. Importantly, they have considerably higher S/N compared to individual spectra, as demonstrated by the comparison of Fig. 8 with Fig. 6. This allows a better characterization of the spectral models, even though some parameters (such as the Comptonization fraction, or the Gaussian line parameters) remain poorly constrained.

Restricting the fit to the 2–10 keV range, to avoid complications with the putative soft excess, the high-flux spectrum appears significantly steeper ($\Gamma = 1.90^{+0.04}_{-0.02}$) than the low-flux spectrum

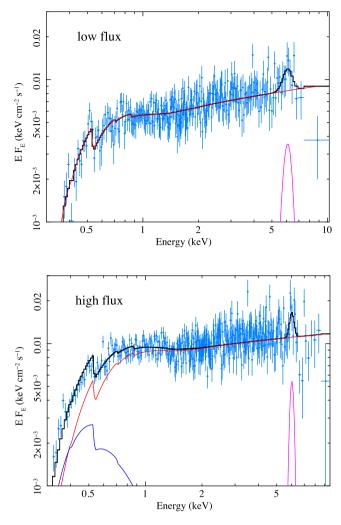


Figure 8. Unfolded (EEUFSPEC command in XSPEC) XRT spectra of Ark 120 for the combined low- and high-flux observations.

 $(\Gamma = 1.72^{+0.07}_{-0.06})$. The difference in slope between the low- and high-flux spectra is seen in Fig. 8. Note that for unfolded spectra (i.e. plots of EF_E versus E, which are equivalent to the νf_{ν} versus ν plots often used in SED studies of AGN), the slope is given by $2 - \Gamma$, which means that the steeper positive slope observed in the low-flux unfolded spectrum (top panel) corresponds to a lower value of Γ compared to the high-flux spectrum (bottom panel).

When the 0.3–10 keV range is considered, both low- and highflux spectra are reasonably well fitted with a coronal BMC model, and both spectra seem to require a Gaussian line ($EW \sim 100-$ 200 eV), whose addition reduces the χ^2 by 5.4 and 6.5 (for three additional parameters) in the low- and high-flux cases, respectively. However, only the high-flux spectrum requires a second BMC model to fit the soft energy range. The best-fitting values of this fitting procedure are summarized in Table 3.

This is confirmed by the flux-selected spectral fitting analysis which shows that the high-flux spectrum is indeed significantly steeper ($\Gamma = 1.90^{+0.04}_{-0.02}$) than the low-flux spectrum ($\Gamma = 1.72^{+0.07}_{-0.06}$), and that only the high-flux spectrum is statistically improved by the addition of a second BMC model (with $kT \sim 0.03$ keV) to fit the low-energy part of the spectrum.

In summary, the spectral analysis based on model fitting of fluxselected spectra confirms the existence of spectral variability during

Table 3. Results from the spectral fitting of high and low flux of Ark 120.

Results	Low flux	High flux
χ^2/dof	1233.5/261	404.6/380
kT_1 (keV)	0.13 ± 0.01	0.15 ± 0.01
α_1	0.80 ± 0.04	0.87 ± 0.04
$\log(A_1)$	0.08	0.2
Norm _{BMC1}	$(2.6 \pm 0.1) \times 10^{-4}$	$(3.5 \pm 0.1) \times 10^{-4}$
kT_2 (keV)	_	0.07 ± 0.014
α2	-	4
$\log(A_2)$	_	0.2
Norm _{BMC2}	-	$(3.5 \pm 0.1) \times 10^{-4}$
E_{line} (keV)	6.3 ± 0.1	6.2 ± 0.1
σ_{line} (keV)	0.35	0.1
Norm _{line}	$(8.5 \pm 2.5) \times 10^{-5}$	$(5.2 \pm 2.5) \times 10^{-5}$

Note. The errors given in this table are 1σ errors.

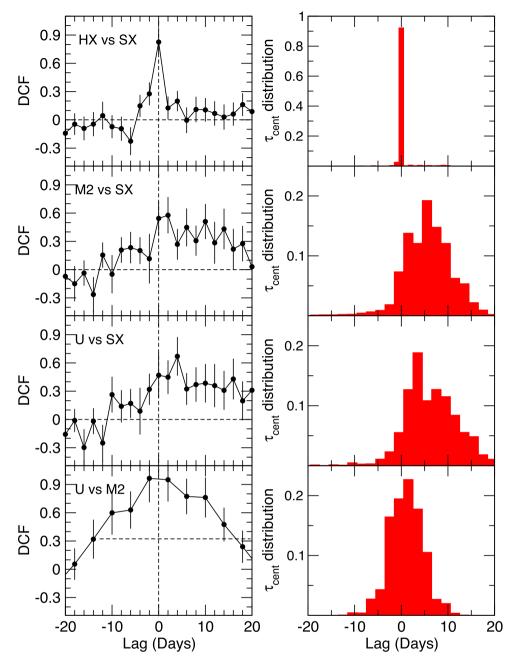
the *Swift* campaign of Ark 120 revealed by the model-independent spectral variability analysis, with steeper spectra observed when the source has higher count rate.

5 CORRELATION ANALYSIS

We used the discrete correlation function (DCF) method of Edelson & Krolik (1988) to compute the correlation function (CCF) at lags $k = 0, \pm l\Delta t$, where $l = 1, ..., 10, \Delta t = 2$ d. As a reference, we used the soft X-ray light curve (hereafter, SX indicates the energy range 0.3–1 keV, and HX the hard X-ray light curve in the 1–10 keV range). We computed the HX versus SX, M2 versus SX, *U* versus SX, and *U* versus M2 cross-correlations. In the case of the *U* versus M2 correlation, we considered the M2 light curve as reference light curve. Positive lags mean that the reference light curve follows. We calculate the centroid of the DCF, τ_{cent} , as the mean of all the DCF points which are >0.75 × DCF_{max}, and we accept it as our estimate of the time lag between two light curves. We also compute the average DCF_{max} as the mean of the DCF values of the same points.

The resulting CCFs are shown in the left-hand panels of Fig. 9. The HX versus SX shows a strong, narrow peak at zero lag. On the other hand, the UV/X-ray correlations are skewed towards positive lags, suggesting that the SX band variations lead those in the UV band. DCF_{max} values are smaller in the cross-correlations between the UV light curves and the X-ray band. This is not surprising, given the fact that the UV band light curves are much 'smoother' than the X-ray band light curves (see Fig. 1). Finally, when U is cross-correlated with the UVM2 band, the CCF is roughly symmetric and shows a strong peak of the order of DCF_{max} ~ 0.9.

In Table 4, we list DCF_{max} together with the time lags, τ_{cent} , between the various bands along with their respective 90 per cent errors. The errors were estimated using the Monte Carlo simulation method proposed by Peterson et al. (1998). For each light curve pair which we cross-correlated, we produced 10 000 simulated light curves following their "random subset selection" prescription. We computed the DCF of each light curve pair, τ_{cent} , and DCF_{max} exactly as we did with the observed light curves. We used the 10 000 values to build up the τ_{cent} and DCF_{max} distribution function. The distribution of the centroid time lags are also plotted in Fig. 9 (right-hand panels). We used these distributions to estimate the 90 per cent confidence limits, which we assume are representative of the 90 per cent confidence limits of the computed τ_{cent} and DCF_{max} values when using the observed light curves.



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Figure 9. Left-hand panels: plots of cross-correlation between the 0.3 and 1 keV soft X-ray flux and the hard (1–10 keV) X-rays (top panel), UVM2 flux (second panel), *U* flux (third panel), and *U* versus UVM2 (bottom panel). Right-hand panels: distributions of the centroid time lags for the various cross-correlations.

We do not find a statistically significant detection of delays between any of the light curve pairs considered: all the "lags" listed in Table 4 are consistent with zero within their 90 per cent confidence limits. In the case of the hard versus soft X–ray, the delays are very small (variations in HX and SX happen almost simultaneously). On the other hand, positive values of τ of the order of a few (~4) days are tentatively detected in the case of the cross-correlation between the UV light curves and the soft X-ray band. The lags are identical when we consider the cross-correlation between the UVM2 and the U-band light curves.

In order to reduce the statistical uncertainty in this crosscorrelation analysis and better constrain the delay between X-ray and UV light curves in Ark 120, we tried to combine the U and UVM2 light curves, which have similar trends and do not show any significant delay. To this end, we first interpolated the *U* light curve, and then shifted the *UVM2* light curve by a multiplicative factor obtained by minimizing the rms between the interpolated *U* values and the shifted *UVM2* values. The resulting combined light curve (hereafter, *U+UVM2*) is simultaneous to the X-ray light curve and has the same number of data points. The CCF analysis with this combined light curve, shown in Fig. 10, reveals that the soft X-ray light curve leads the UV one by 7.5 \pm 7 d (errors indicate the 90 per cent confidence limits). The results are not affected by the uncertainty in the scaling factor when we create the combined *U+UVM2* light curve, due to the small uncertainty in the *F*_{var} of the individual light curves. For completeness, we also performed a

Energy bands (1)	τ_{cent} (d) (2)	DCF _{max} (3)
HX versus SX	0 ± 0.5	$0.83^{+0.17}_{-0.13}$
UVM2 versus SX	$4.5^{+9.5}_{-6.5}$	$0.52_{-0.02}^{+0.45}$
U versus SX	$4^{+12}_{-5.5}$	$0.67^{+0.30}_{-0.08}$
U versus UVM2	4^{+3}_{-8}	$0.86^{+0.15}_{-0.07}$
(U+UVM2) versus SX	7.5 ± 7	$0.52^{+0.24}_{-0.12}$
HR versus SX	0^{+1}_{-10}	$-0.79^{+0.26}_{-0.21}$
HR versus (U+UVM2)	-15^{+8}_{-6}	$-0.48\substack{+0.39\\-0.08}$

Notes. Columns Table 4: 1 = correlated light curves. 2 = lags measured in days with the 90 per cent errors. 3 = maximum of DCF with the 90 per cent errors.

cross-correlation between the SX and the X-ray hardness ratio, HR, and between the combined UV light curve and HR. In both cases, the flux light curves were the reference light curves. The results of this analysis indicate that HR is anticorrelated with both SX and (U+UVM2).

In summary, our cross-correlation analysis confirms that soft and hard X-ray variations are strongly correlated and occur nearly simultaneously. Similarly, the variations in the *U* and *UVM2* filters appear to be correlated and without substantial delay. When the UV light curves are correlated with the X-ray light curve, a possible (but not statistically significant) delay is suggested, with the X-ray leading the variations in the UV bands by a few days. This result is confirmed at a higher significance level when the *U* and *UVM2* are combined and then correlated with the soft X-ray.

6 SUMMARY AND CONCLUSION

We first summarize the most relevant results of our *Swift* monitoring campaign of Ark 120, and then discuss their implications in the broader context of AGN variability studies.

(i) Temporal variability – the six-month XRT and UVOT monitoring of Ark 120 revealed that strong variability in the X-ray and UV bands, observed in past pointing observations on shorter timescales, occurs on all time-scales probed by the *Swift* campaign, i.e. from a few days to a few months (see Figs 1 and 2). While the X-ray variability is characterized by frequent large-amplitude changes where the count rate can double or halve in periods as short as 2–4 d, the variations observed in the UV bands are smoother with flux changes of the order of 20–40 per cent occurring on time-scales of months. This different behaviour can be quantified by fractional variability measurements: F_{var} increases from ~5 per cent in the *U* band, to ~8 per cent in *UVM2*, up to ~20 per cent in the 0.3–10 keV energy band.

(ii) Spectral variability – the continuous temporal variability of Ark 120 appears to be associated with persistent spectral variability based on various model-independent analyses. For example, the light curve of the hardness ratio HR = (h - s)/(h + s) is inconsistent with the hypothesis of constancy at a high confidence level. When HR is plotted versus the total X-ray count rate, a weak but statistically significant anticorrelation is found, indicating that the X-ray spectrum softens when the source brightens, which is the typical behaviour observed in Seyfert galaxies (see Fig. 4). Additionally, soft and hard X-ray fluxes are tightly correlated and well described by a linear equation, whose slope is consistent with unity and whose intercept is inconsistent with zero, suggesting the presence of a constant hard component. Finally, the light curve of the broad-band spectral index indicates that the entire SED varies throughout the monitoring campaign (see Fig. 5).

(iii) Spectral analysis – the spectral analysis of individual ≤ 1 ks observations does not provide conclusive results about the long-term spectral variability of Ark 120, due to the limited statistics. However, combining several individual spectra into a low-flux and a high-flux spectrum and then performing a model fitting of these two flux-selected spectra makes it possible to conclude that the steeper-when-brighter behaviour is caused by the steepening of the photon index.

(iv) Correlation analysis – a cross-correlation analysis of the *Swift* UVOT and XRT light curves of Ark 120 indicates that soft and hard X-ray variations are strongly correlated and occur nearly simultaneously. Also the *U* and *UVM2* light curves are well correlated with each other and do not show any substantial delay. When the UV light curves are correlated with the X-ray light curve, a possible (but not statistically significant) delay is tentatively detected,

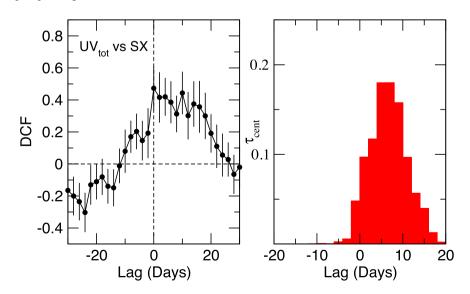


Figure 10. Left-hand panel: plots of cross-correlation between the combined (*U*+UVM2) light curve and the soft X-ray energy band 0.3–1 keV. Right-hand panel: distribution of the centroid time lags for this cross-correlation.

with the X-ray leading the variations in the UV bands by a few days (see Fig. 9). This result is confirmed at a higher significance level by using the combined U and UVM2 light curve for the correlation analysis with the soft X-ray light curve (see Fig. 10). Finally, the hardness ratio HR appears to be anticorrelated with the soft X-ray light curve, and with the combined UV light curve. While there is no relevant lag between HR and X-ray flux, it appears that the changes in the UV light curve are delayed by several days ($\tau = -15^{+8}_{-6}$ d) with respect to the HR changes. The reason for the UV-HR correlation was to investigate the possibility that the observed UV photons are the input soft photons up-scattered in the hot corona. In this case, an increase in the flux of soft photons may cause the cooling of the corona, and hence a steepening in the observed X-ray spectrum, as it has been observed in the past (e.g. Nandra et al. 2000). Our results do not support this possibility and may be explained by the fact that the X-rays and HR are strongly anticorrelated (with no delay), and the X-rays and UV are moderately correlated with a delay of a few days.

Our study confirms that Ark 120 behaves as a typical Seyfert galaxy with persistent X-ray (and UV) flux variability associated with spectral variability, where the spectrum softens as the source brightens. This spectral behaviour too is common to the vast majority of Seyfert galaxies. However, the 'bare' nature of Ark 120 ensures that the spectral variability is caused by intrinsic changes in the primary emission, rather than being associated with variations of the absorber surrounding the source, as suggested for many other AGN. Our results, obtained from a model-independent analysis of monitoring data spanning several months, appear to be consistent with those based on detailed spectral analysis of broad-band spectra obtained from the long uninterrupted exposures. For instance, the fact that the intercept of the hard versus soft X-ray correlation is positive and inconsistent with zero indicates the presence of a constant hard component, which is naturally explained by the reflection component detected by Matt et al. (2014) using high-quality spectra from XMM-Newton and NuSTAR.

The results from the cross-correlation analysis (the tentative time lag of the UV flux with respect to the X-ray light curve) are consistent with the reprocessing scenario, where changes in the UV/optical-emitting accretion disc are driven by changes in the X-ray corona. Although the measured time lag is poorly constrained due to the large statistical uncertainty ($\tau = 7.5 \pm 7$ d, which is obtained when the combined UV light curve is used for the correlation analysis), to put Ark 120 in context, it is helpful to compare its correlation results with those obtained in similar studies.

In particular, it is interesting to investigate whether there exists a correlation between time lags and $M_{\rm BH}$ or \dot{m} , using AGN whose UV and X-ray light curves have been simultaneously monitored by *Swift* for several months. In addition to the two objects studied by our group: PKS 0558–504 – $M_{\rm BH} \sim 3 \times 10^8 \, {\rm M}_{\odot}$ and $\dot{m} \ge 1-$ (Gliozzi et al. 2013), and Ark 120 – $M_{\rm BH} = 1.5 \times 10^8 \, {\rm M}_{\odot}$ and $\dot{m} = 0.005-$ (this work), this limited sample of AGN comprises NGC 4395 – $M_{\rm BH} = 3.6 \times 10^5 \, {\rm M}_{\odot}$ and $\dot{m} = 0.005-$ (Cameron et al. 2012), NGC 2617 – $M_{\rm BH} = 4 \times 10^7 \, {\rm M}_{\odot}$ and $\dot{m} = 0.03-$ (Edelson et al. 2015), and NGC 6814 – $M_{\rm BH} = 2.6 \times 10^6 \, {\rm M}_{\odot}$ and $\dot{m} = 0.01-$ (Pancoast et al. 2014; Troyer et al. 2016).

In Fig. 11, we plotted the time delays between the UV band and the X-ray detected for these objects. For all objects, we used the lag value reported for the UVOT U filter ($\lambda_{\text{peak}} = 350 \text{ nm}$), with the exception of NGC 6814 for which only the lag of UVW1 ($\lambda_{\text{peak}} = 260 \text{ nm}$) was measured. PKS 0558–504 was not included

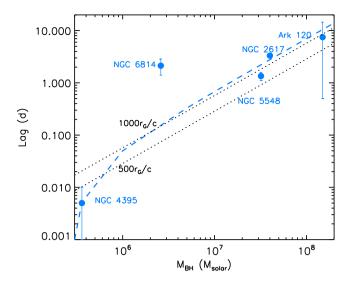


Figure 11. Time lags of the UV band with respect to the X-ray band plotted versus the BH mass. The dashed line represents the best-fitting linear model UV/X-ray⁻Lag = $-0.02 \pm 0.01 + (7.1 \pm 0.4) \times 10^{-8} M_{BH}$. The dotted dark lines represent the light travel times for 500 r_G/c (bottom) and 1000 r_G/c (top).

because the detection of a lag was not statistically significant. However, it is worth noting that PKS 0558–504 putative delay $(\tau = -16.8^{+16.8}_{-14.7} \text{ d})$ was of the same order as the one detected in Ark 120, which has a similar BH mass, but negative (i.e. with the UV emission which appears to lead the X-rays by a few days). It is interesting to note that PKS 0558–504 is the only object for which such a negative delay has been putatively detected and the only AGN of this sample with accretion rate above the Eddington level.

All objects in Fig. 11, with the exception of NGC 6814, are reasonably well fitted with a linear model, represented by the dashed line UV/X-ray Lag = $-0.02 \pm 0.01 + (7.1 \pm 0.4) \times 10^{-8} M_{\rm BH}$. This finding is qualitatively in agreement with the general picture of BH systems, where the length-scale is set by the BH mass, naturally implying a larger physical separation (and hence longer delays) for systems with larger $M_{\rm BH}$.

The dotted lines, which represent the light travel times for 500 r_G/c and 1000 r_G/c , suggest that, with the exception of NGC 4395 (for which no significant lag was detected), all AGN of this sample require a physical separation between the X-ray-emitting region and the UV region of the order of 1000 r_G or more (NGC 6814). These values are considerably larger than the physical locations of the UV-emitting region predicted by the standard accretion disc model; using equation 2 from Cameron et al. (2012) we obtain values of the order $100-250r_G$. We therefore conclude that, for this sample of AGN, these cross-correlation results imply a larger accretion disc compared to the Shakura–Sunyaev standard model, as suggested by recent findings based on micro-lensing studies from Mosquera et al. (2013) and intensive simultaneous monitoring of several energy bands in NGC 5548 (Edelson et al. 2015).

In conclusion, our work indicates that long-term monitoring studies of AGN provide useful information which is complementary to that obtained in long-exposure spectral studies. Importantly, combining Ark 120 correlation results with those of similar studies of AGN monitored by *Swift*, suggests the existence of a positive correlation between time lags and BH mass. Additional monitoring studies of AGN spanning a broader range of $M_{\rm BH}$ and M_{\odot} are necessary to derive a firmer conclusion.

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SUPPORTING INFORMATION

Additional Supporting Information may be found in the online version of this article:

Table 1. Observation log of Ark 120.

Table 2. *Swift* XRT count rates and HR and UVOT fluxes (mJy) of Ark 120.

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